



A multiscale method for linear heterogeneous poroelasticity

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Outline

- ◇ Linear Poroelasticity
- ◇ Classical FEM
- ◇ Multiscale Methods
- ◇ Practical Aspects
- ◇ Numerical Example
- ◇ Outlook: Semi-Explicit Time Discretization
- ◇ Conclusion

Linear Poroelasticity

- ◊ Deformation of porous media (heterogeneous) saturated by an incompressible viscous fluid
- ◊ Couples fluid flow with behavior of surrounding solid
- ◊ Quasi-static: internal equilibrium preserved at any time
- ◊ Applications:
 - ▷ geomechanical modeling (reservoir engineering)
 - ▷ human anatomy for medical applications
 - ▷ similar to linear thermoelasticity
- ◊ Variables:
 - ▷ displacement u , pressure p
 - ▷ both averaged across (infinitesimal) cubic elements
 - ▷ both treated as variables on the entire domain
 - 3 scales!

Linear Poroelasticity

- ◇ On a Lipschitz domain $D \subset \mathbb{R}^d$,

$$\begin{aligned}-\nabla \cdot (\sigma(u)) + \nabla(\alpha p) &= 0, \\ \partial_t \left(\alpha \nabla \cdot u + \frac{1}{M} p \right) - \nabla \cdot \left(\frac{\kappa}{\nu} \nabla p \right) &= f\end{aligned}$$

- ◇ Initial condition for p , boundary conditions
- ◇ Heterogeneous media \rightarrow oscillatory $\mu, \lambda, \kappa, \alpha$

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- ◇ Initial condition for p , boundary conditions
- ◇ Heterogeneous media \rightarrow oscillatory $\mu, \lambda, \kappa, \alpha$
- ◇ Variational form:

$$\begin{aligned}\mathbf{a}(u, v) - \mathbf{d}(v, p) &= 0, \\ \mathbf{d}(\partial_t u, q) + c(\partial_t p, q) + \mathbf{b}(p, q) &= (f, q),\end{aligned}$$

for all $v \in V := [H_0^1(D)]^d$, $q \in Q := H_0^1(D)$

Linear Poroelasticity

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- ◇ Initial condition for p , boundary conditions
- ◇ Heterogeneous media \rightarrow oscillatory $\mu, \lambda, \kappa, \alpha$
- ◇ Bilinear forms $a: V \times V \rightarrow \mathbb{R}$, $b, c: Q \times Q \rightarrow \mathbb{R}$, and $d: V \times Q \rightarrow \mathbb{R}$ defined by

$$\begin{aligned}a(u, v) &:= \int_D \sigma(u) : \varepsilon(v) \, dx, & b(p, q) &:= \int_D \frac{\kappa}{\nu} \nabla p \cdot \nabla q \, dx, \\ c(p, q) &:= \int_D \frac{1}{M} pq \, dx, & d(u, q) &:= \int_D \alpha (\nabla \cdot u) q \, dx\end{aligned}$$

Classical FEM

- ◇ V_h, Q_h : piecewise affine finite element spaces on \mathcal{T}_h
- ◇ Backward Euler in time with uniform time steps τ
- ◇ Discrete derivative $D_\tau u_h^n := (u_h^n - u_h^{n-1})/\tau$
- ◇ For $n \in \{1, \dots, N\}$ find $u_h^n \in V_h$ and $p_h^n \in Q_h$ such that

$$\begin{aligned} a(u_h^n, v) - d(v, p_h^n) &= 0, \\ d(D_\tau u_h^n, q) + c(D_\tau p_h^n, q) + b(p_h^n, q) &= (f^n, q) \end{aligned}$$

for all $v \in V_h$ and $q \in Q_h$

- ◇ Properties:
 - ▷ unique solution at each time step
 - ▷ stability estimates as in [MalP17]

Classical FEM

Theorem ([ErnM09])

For coefficients $\mu, \lambda, \kappa, \alpha \in W^{1,\infty}(D)$ the error of the exact solution (u, p) and the fully discrete solution (u_h^n, p_h^n) is bounded by

$$\|u(t_n) - u_h^n\|_1 + \|p(t_n) - p_h^n\| \leq C_\epsilon h + C\tau,$$

where the constant C_ϵ scales with the maximum of $\|\mu\|_{W^{1,\infty}(D)}$, $\|\lambda\|_{W^{1,\infty}(D)}$, $\|\kappa\|_{W^{1,\infty}(D)}$, and $\|\alpha\|_{W^{1,\infty}(D)}$.

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- ◊ Expected convergence rate $h + \tau$
- ◊ Involved constant C_ϵ scales with ϵ^{-1} in oscillatory media with period ϵ
→ unfeasible approach

Multiscale Spaces and Projections

- ◇ Larger mesh size $H > h$, \mathcal{T}_h refinement of \mathcal{T}_H
- ◇ Defines $V_H \subseteq V_h$ and $Q_H \subseteq Q_h$
- ◇ Goals:
 - ▷ construct **new discrete function space** with the same dimension as $V_H \times Q_H$
 - ▷ **better approximation properties**
 - ▷ *localized orthogonal decomposition* (LOD) [MaIP14]

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- ◇ **Quasi-interpolation** operator $I_H: V_h \rightarrow V_H$

$$H_K^{-1} \|v - I_H v\|_{L^2(K)} + \|\nabla I_H v\|_{L^2(K)} \lesssim \|\nabla v\|_{L^2(\omega_K)}$$

- ▷ example: **local L^2 -projection + averaging** in nodes
- ▷ analogously $I_H: Q_h \rightarrow Q_H$

Multiscale Spaces and Projections

- ◊ Fine scale spaces

$$V_{\text{fs}} := \{v \in V_h : I_H v = 0\}, \quad Q_{\text{fs}} := \{q \in Q_h : I_H q = 0\}$$

- ◊ Leads to the decomposition

$$V_h = V_H \oplus V_{\text{fs}}, \quad Q_h = Q_H \oplus Q_{\text{fs}}$$

- ◊ Correctors $\mathcal{C}_{\text{fs}}^1 : V_h \rightarrow V_{\text{fs}}$, $\mathcal{C}_{\text{fs}}^2 : Q_h \rightarrow Q_{\text{fs}}$

$$a(\mathcal{C}_{\text{fs}}^1 u, v) = a(u, v), \quad b(\mathcal{C}_{\text{fs}}^2 p, q) = b(p, q), \quad v \in V_{\text{fs}}, \quad q \in Q_{\text{fs}}$$

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- ◇ New finite element spaces

$$V_{\text{ms}} := \{v_H - \mathcal{C}_{\text{fs}}^1 v_H : v_H \in V_H\},$$

$$Q_{\text{ms}} := \{q_H - \mathcal{C}_{\text{fs}}^2 q_H : q_H \in Q_H\}$$

and *a/b*-orthogonal decompositions

$$V_h = V_{\text{ms}} \oplus V_{\text{fs}}, \quad Q_h = Q_{\text{ms}} \oplus Q_{\text{fs}}$$

First Multiscale Method

- ◇ Multiscale method from thermoelasticity [MalP17]
- ◇ Idea: consider **stationary system**
 - solve time-independent corrector problem
- ◇ Problems:
 - ▷ two **coupled correctors**
 - ▷ need additional fine scale correction
- ◇ Find $\tilde{u}_{\text{ms}}^n = u_{\text{ms}}^n + u_{\text{fs}}^n \in V_{\text{ms}} \oplus V_{\text{fs}}$, $\tilde{p}_{\text{ms}}^n \in Q_{\text{ms}}$ such that

$$a(\tilde{u}_{\text{ms}}^n, v) - d(v, \tilde{p}_{\text{ms}}^n) = 0,$$

$$d(D_\tau \tilde{u}_{\text{ms}}^n, q) + c(D_\tau \tilde{p}_{\text{ms}}^n, q) + b(\tilde{p}_{\text{ms}}^n, q) = (f^n, q),$$

$$a(u_{\text{fs}}^n, w) + d(w, \tilde{p}_{\text{ms}}^n) = 0$$

for $v \in V_{\text{ms}}$, $q \in Q_{\text{ms}}$, and $w \in V_{\text{fs}}$

Alternative Multiscale Method

- ◇ New approach: exploit saddle point structure
- ◇ Advantages:
 - ▷ decoupled corrector problems
 - ▷ no additional fine scale correction
 - ▷ independent of the bilinear form d (and thus α)
- ◇ Find $u_{\text{ms}}^n \in V_{\text{ms}}$ and $p_{\text{ms}}^n \in Q_{\text{ms}}$ such that

$$a(u_{\text{ms}}^n, v) - d(v, p_{\text{ms}}^n) = 0,$$

$$d(D_\tau u_{\text{ms}}^n, q) + c(D_\tau p_{\text{ms}}^n, q) + b(p_{\text{ms}}^n, q) = (f^n, q)$$

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$$\begin{aligned} a(u_{\text{ms}}^n, v) - d(v, p_{\text{ms}}^n) &= 0, \\ d(D_\tau u_{\text{ms}}^n, q) + c(D_\tau p_{\text{ms}}^n, q) + b(p_{\text{ms}}^n, q) &= (f^n, q) \end{aligned}$$

Theorem

For consistent initial data $u_h^0 \in V_h$, $p_h^0 \in Q_h$ and $u_{\text{ms}}^0 \in V_{\text{ms}}$, $p_{\text{ms}}^0 \in Q_{\text{ms}}$, the error of the multiscale solution compared to the fine scale solution satisfies

$$\|u_h^n - u_{\text{ms}}^n\|_1 + \|p_h^n - p_{\text{ms}}^n\|_1 \lesssim H \text{data}^n + t_n^{-1/2} H \|p_h^0\|_1$$

Practical Aspects

- ◆ corrector problems are only computed for basis $\{\Lambda_j\}_{j=1}^M$ of V_H :

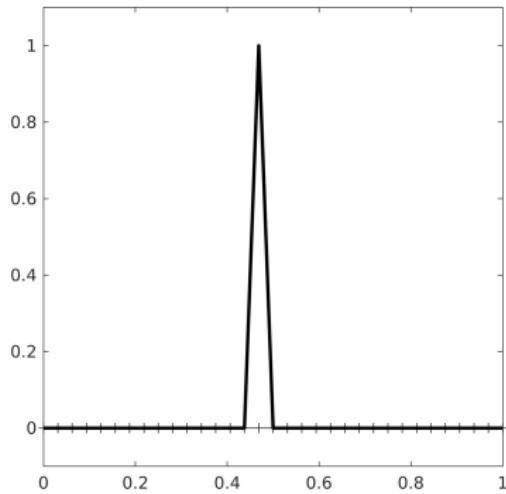
$$a(\mathcal{C}_{fs}^1 \Lambda_j, w) = a(\Lambda_j, w)$$

for all $w \in V_{fs}$

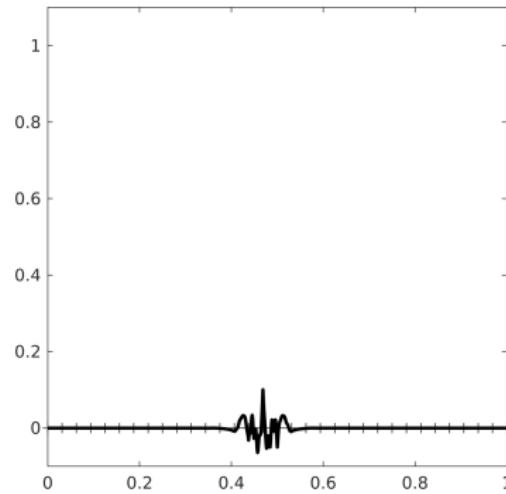
→ support possibly **global**!

- ◆ problems are **localized** on patches
→ same error rate with patch size $|H \log H|$

Illustration in 1D

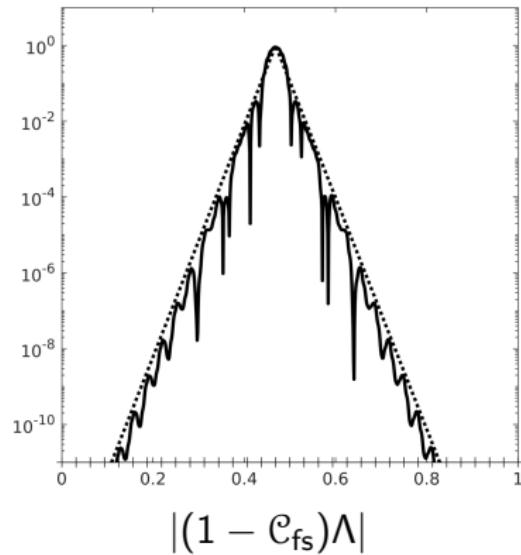
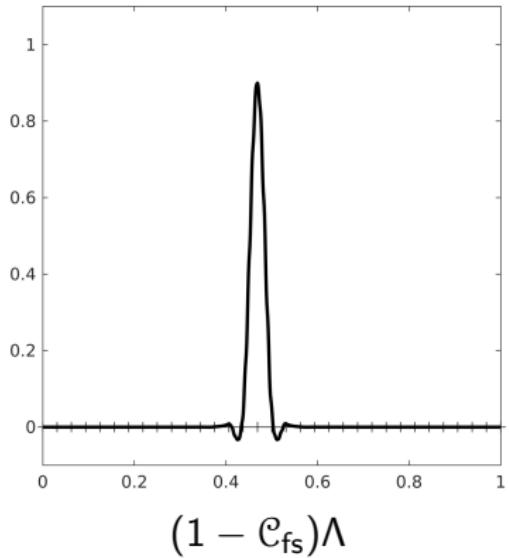


Hat function Λ



Correction $C_{fs}\Lambda$

Illustration in 1D



Numerical Example

- ◇ Parameters

$$D = (0, 1)^2, \quad T = 1, \quad \tau = 0.01, \quad h = \sqrt{2} \cdot 2^{-8} \approx 0.0055$$

- ◇ Reference solution on uniform fine grid \mathcal{T}_h

- ◇ Error measure

$$\|(\boldsymbol{v}, q)\|_{D,N}^2 := \sum_{i=1}^N \tau \left(\|\nabla \boldsymbol{v}^i\|^2 + \|\nabla q^i\|^2 \right)$$

- ◇ Piecewise constant parameters on \mathcal{T}_ϵ ,

$$\mu \sim U[32.2, 62.2], \quad \lambda \sim U[41, 61],$$

$$\kappa \sim U[0.1, 0.12], \quad \alpha \sim U[0.5, 1]$$

Numerical Example

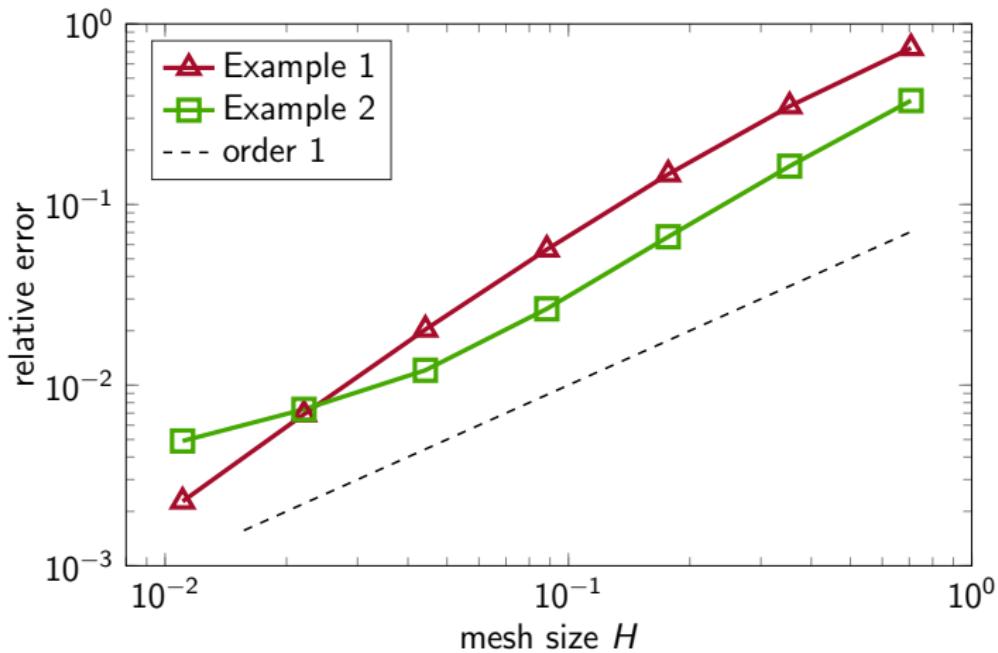
◇ Example 1

- ▷ $f = 1$
- ▷ initial data $p^0(x) = (1 - x_1)x_1(1 - x_2)x_2$
- ▷ hom. Dirichlet boundary conditions for p
- ▷ mixed boundary conditions for u

◇ Example 2

- ▷ random f with values between 0 and 1
- ▷ initial data $p^0(x) = (1 - x_2)x_2$
- ▷ hom. Dirichlet boundary conditions on the top,
hom. Neumann boundary conditions otherwise

Numerical Example



Outlook: Semi-Explicit Scheme

Implicit discretization:

$$\begin{aligned} a(u_h^n, v) - d(v, p_h^n) &= 0, \\ d(D_\tau u_h^n, q) + c(D_\tau p_h^n, q) + b(p_h^n, q) &= (f^n, q) \end{aligned}$$

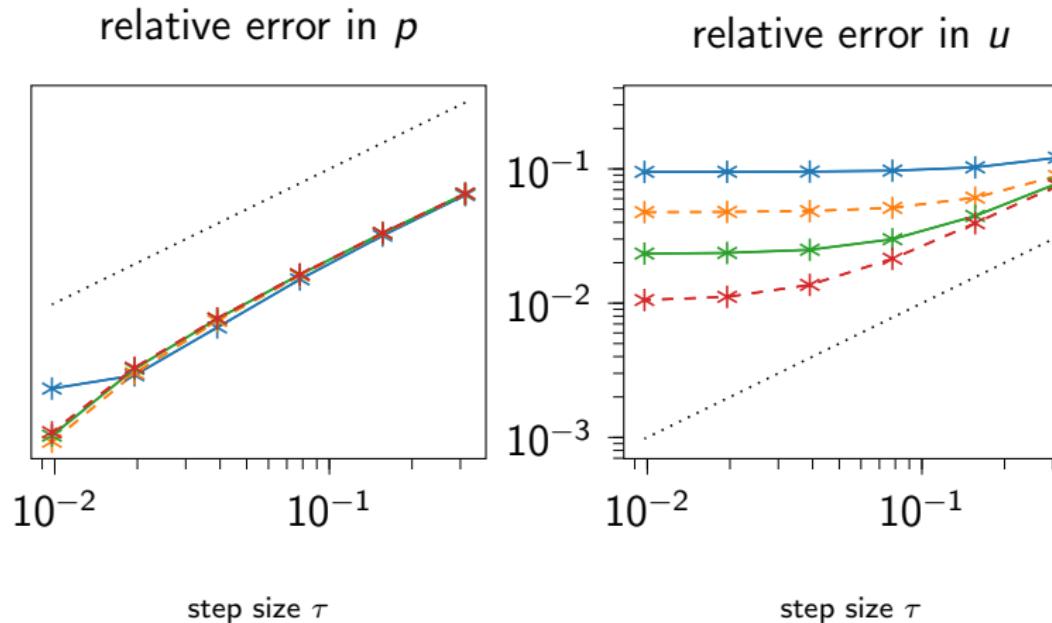
Semi-Explicit discretization under weak coupling [AltMU19]:

$$\begin{aligned} a(u_h^n, v) - d(v, p_h^{n-1}) &= 0, \\ d(D_\tau u_h^n, q) + c(D_\tau p_h^n, q) + b(p_h^n, q) &= (f^n, q) \end{aligned}$$

Benefit: decoupled equations, faster computations

→ combine this with spatial numerical homogenization!

Illustration: Energy Errors



lines $\hat{=}$ semi-explicit, stars $\hat{=}$ implicit scheme, $h = 2^{-\{5,6,7,8\}}$

Conclusion

- ◇ Linear heterogeneous poroelasticity: 3 scales
- ◇ Standard FEM approach unfeasible
- ◇ New multiscale method
- ◇ Decoupled corrector problems, independent of α
- ◇ Practical implementations include localization
- ◇ Optimal first-order convergence theoretically and numerically

Future work: semi-explicit schemes

THANK YOU FOR YOUR ATTENTION!

<http://scicomp.math.uni-augsburg.de>

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